



# Control of metamagnetic phase transition in epitaxial FeRh films by changing atomic order degree

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## ABSTRACT

Near-equiatomic FeRh alloys undergo a first-order metamagnetic phase transition around room temperature. Here, we show that the phase transition parameters of epitaxial FeRh films can be adjusted to a great extent through changing the growth temperature, annealing temperature, and annealing time. The influences of the surface morphology and Fe (Rh) element distribution before and after heat treatments have been analyzed. The experimental results show that the increase of annealing temperature reduces the displacement of Fe (Rh) atoms. As a result, the resistivity is reduced, the phase transition interval is narrowed, and the phase transition is shifted to high temperature. The experimental results indicate that the first-order phase transition parameters in epitaxial FeRh films can be controlled by heat treatments and are correlated with the atomic order degree of Fe (Rh).

## 1. Introduction

Near-equiatomic FeRh alloys exhibit metamagnetic transition from antiferromagnetic (AF) to ferromagnetic (FM) transition around room temperature [1,2], which is accompanied by changes in magnetization [3,4], entropy [5], lattice parameter [6], and resistivity [7]. Fig. 1(a) shows a typical temperature dependent magnetization curve of first-order metamagnetic transition with several important parameters: critical AF-FM transition temperature ( $T_t$ ), transition temperature range ( $W_{AF-FM}$ ), and thermal hysteresis width ( $W_{TH}$ ) [8]. From the application point of view,  $T_t$  determines the operating temperature [9], while  $W_{AF-FM}$  and  $W_{TH}$  are critical to the performance. For example, in the application of magnetic refrigeration technology [10], it needs a narrow  $W_{TH}$  and  $W_{AF-FM}$  to obtain a greater giant magnetocaloric effect. For the successful production of controllable FM pattern in a non-magnetic background [11], the  $W_{TH}$  of the FeRh films must be wide enough to resist thermal disturbance and narrow enough for effective control. The  $T_t$  is near room temperature to facilitate the observation of the FM pattern. For the construction of some spintronic devices [12,13], a wide

$W_{TH}$  or  $W_{AF-FM}$  is required to ensure a stable coexistence state with both AF and FM phases in a suitable temperature range. Therefore, the ability to control the first-order magnetic phase transition parameters ( $T_t$ ,  $W_{AF-FM}$  and  $W_{TH}$ ) of FeRh films is particularly important for the applications.

In the past, people mainly focused on the control of  $T_t$  through different methods, such as magnetic field [14], composition adjustment [15,16], electric field induced strain [17], chemical displacement [18], and ion irradiation [19], and heat treatment [20–22]. For epitaxial FeRh films, the ion irradiation can also cause the surface layer of the film to be FM [23], while the substrate/film interface is in the AF state. Composition adjustment can produce jump-type phase transition [15]. There are few reports on the effect of heat treatment on the first-order magnetic phase transition of epitaxial FeRh films [24]. Furthermore, comparing FeRh bulk and film samples, it is found that as the annealing temperature increases, the  $T_t$  of the bulk sample decreases [25] while the film samples show the opposite properties [24]. Therefore, a systematic study of the heat treatment process is very important for the regulation of the first-order magnetic phase transition in epitaxial FeRh films. In addition, the control mechanism of heat treatment on the first-

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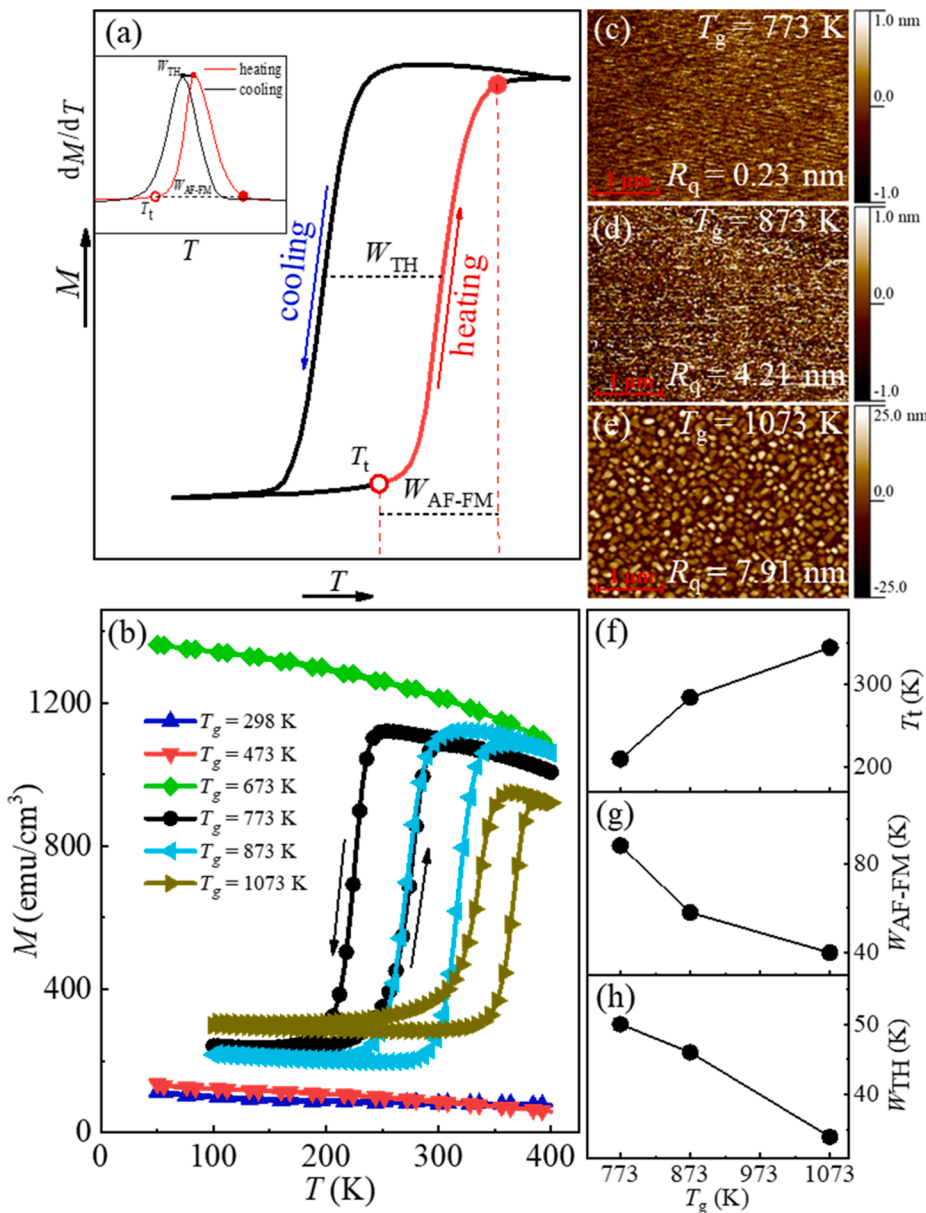
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**Fig. 1.** (a) The temperature-dependent magnetization curve diagram is marked with first-order magnetic phase transition parameters ( $T_t$ ,  $W_{AF-FM}$ ,  $W_{TH}$ ). The inset shows the relationship between  $T$  and  $dM/dT$ . The locations where  $T_t$ ,  $W_{AF-FM}$  and  $W_{TH}$  are located are marked on the illustration. Where  $T_t$  corresponds to the critical AF-FM transition temperature of the heating curve.  $W_{AF-FM}$  is difference between  $T_t$  in the heating curve and the temperature value corresponding to just reaching saturation magnetization intensity.  $W_{TH}$  is the temperature difference between the highest points of the  $dM/dT$  curve of the heating and cooling curves. (b) Temperature-dependence magnetization curves of FeRh films for different  $T_g$ . Atomic force micrographs were taken at growth temperatures (c)  $T_g = 773$  K, (d)  $T_g = 873$  K and (e)  $T_g = 1073$  K. The scanning size of atomic force micrographs is  $4 \times 3 \mu\text{m}^2$ . (f), (g) and (h) are the magnetic phase transition parameters  $T_t$ ,  $W_{AF-FM}$  and  $W_{TH}$  with  $T_g$ , respectively.

order magnetic phase transition is still under investigation [25].

In this work, we made a systematic investigation of the effect of the thermal treatment on the first-order magnetic phase transition parameters in epitaxial FeRh films. The ways of heat treatment include growth temperature ( $T_g$ ), annealing time ( $t_a$ ), and annealing temperature ( $T_a$ ). All of the heat treatments can regulate the first-order magnetic phase transition parameters of FeRh films. The morphology and Fe (Rh) element distribution of the films before and after heat treatments has been characterized. The experimental results indicate that the first-order magnetic phase transition parameters of the epitaxial FeRh films are correlated with the order degree of Fe (Rh) atoms.

## 2. Experimental

The FeRh films were prepared by an ultra-high vacuum magnetron sputtering system. Before growth, commercial (001)-oriented MgO substrates were annealed at 773 K for 1 h in a chamber with a vacuum degree of below  $1 \times 10^{-8}$  torr. The thicknesses of the FeRh films are about 30 nm, which were controlled by the growth time and were calibrated by x-ray reflectivity. In the present experiments, we used four

kinds of thermal treatments: (1) FeRh films were grown at different  $T_g$ . (2) FeRh films were grown at 773 K, and then annealed at 773 K with different  $t_a$ , respectively. (3) FeRh films were grown at 773 K, and then annealed at different  $T_a$  for 1.5 h, respectively. (4) a FeRh thin film was grown at room temperature, then continuously annealed it under different  $T_a$  for 1.5 h. The thermal treatment processes of all FeRh films were completed in the chamber with a vacuum degree of below  $1 \times 10^{-8}$  torr. After the thermal treatments were completed, the films were cooled to room temperature at 0.4 K/s and then taken out from the vacuum chamber. The composition of the epitaxial FeRh thin film is close to  $\text{Fe}_{51}\text{Rh}_{49}$ , which is verified by the energy dispersive spectrometer.

The atomic force micrographs were collected in tapping mode using a commercial cantilever probe with a resonance frequency of about 75 kHz. The resolution is  $256 \times 256$  pixels, the scan frequency is 1 Hz, and the scan size is  $5 \times 5 \mu\text{m}^2$ . The raw data processing (background subtraction, flattening and filtering) and subsequent atomic force micrograph analysis were performed using Nano Scope Analysis software.

The electrical transport of FeRh films were studied in the physical property measurement system (PPMS). The temperature-dependent electrical transport measurements were performed using the van der

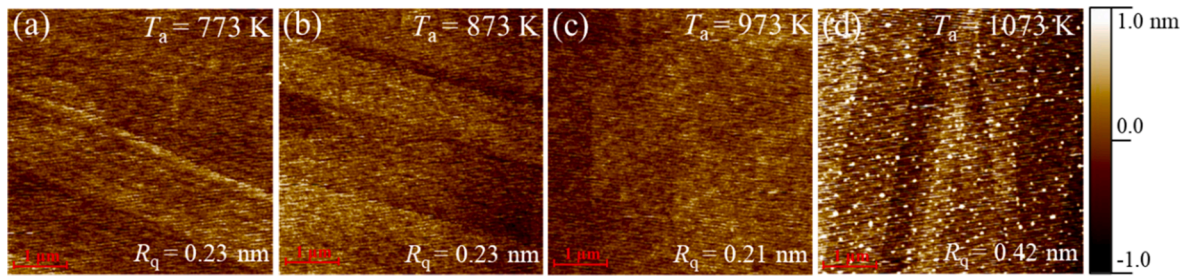


Fig. 2. Atomic force micrographs taken on FeRh films for different  $T_a$ , (a)  $T_a = 773$  K, (b)  $T_a = 873$  K, (c)  $T_a = 973$  K, (d)  $T_a = 1073$  K. The scanning size is  $5 \times 5 \mu\text{m}^2$ .

Pauw geometry with pressed Ag contacts by incrementally cycling the temperature in 2 K steps in the range of 10–400 K. The data of the magnetization in the cycle temperature of 10–400 K were collected with a superconducting quantum interferometer (SQUID-VSM) under a fixed magnetic field of 2 kOe. The temperature was swept at a typical rate of 10 K/min.

The x-ray diffraction (XRD)  $\theta$ - $2\theta$  scans were performed on the D8 Bruker Advance x-ray diffractometer. During the test, a Cu target was used as the incident x-ray source, the incident wavelength was 1.5406 Å, the voltage was 40 kV, the current was 40 mA, the measured angle range was  $20^\circ$ - $70^\circ$ , and the step length of the scanning surface was  $0.01^\circ$ .

The cross-sectional transmission electron microscope (TEM) sample of the 30-nm-thick FeRh films was prepared from a bulk substrate, transferred to a copper grid by a manipulator and ion-milled using an Auriga focused ion-beam (FIB) instrument. Standard  $\text{Ga}^+$  FIB operating procedures were followed (including the final polishing step, adding Pt to protect the electron beam deposition layer and reducing the thickness

of the damaged part) to produce an electronically transparent FeRh/MgO thin layer. The TEM imaging described in this article was performed on a 200 kV full atomic resolution microscope (Talos F200x). A small beam flow was used to reduce the damage during TEM sample preparation. Conventional and high-resolution large-angle circular dark field (HAADF) scanning TEM (STEM) is performed on a cross-sectional transmission electron microscope slice. Mapping provides chemical element distribution analysis.

### 3. Results and discussions

Fig. 1(b) shows the temperature dependent magnetization curves of FeRh films growth at different  $T_g$ . The films display no first-order metamagnetic phase transition when the  $T_g$  is below 673 K, it exhibits either weak-magnetic state (grown at temperature below 473 K) or a FM state (grown at 673 K). The first-order magnetic phase transition begins to appear when the  $T_g$  is higher than 773 K. As the  $T_g$  increases, the  $T_t$

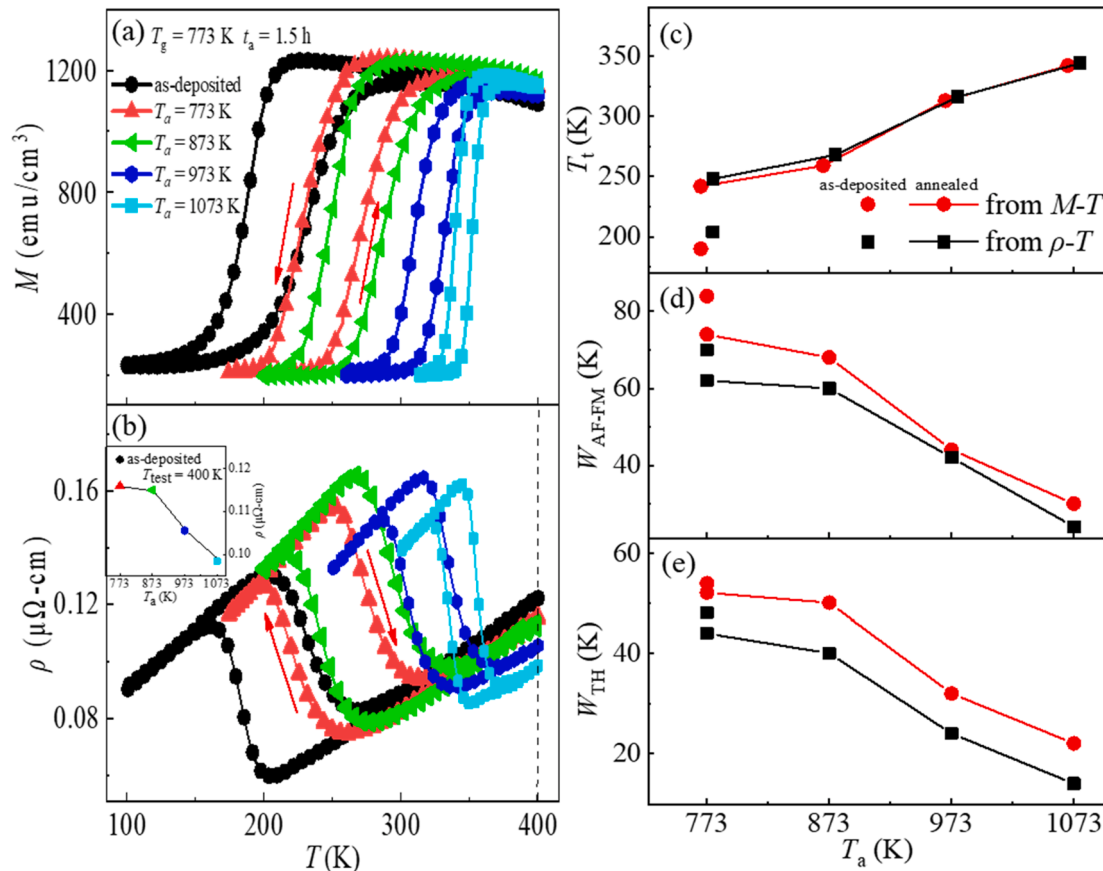
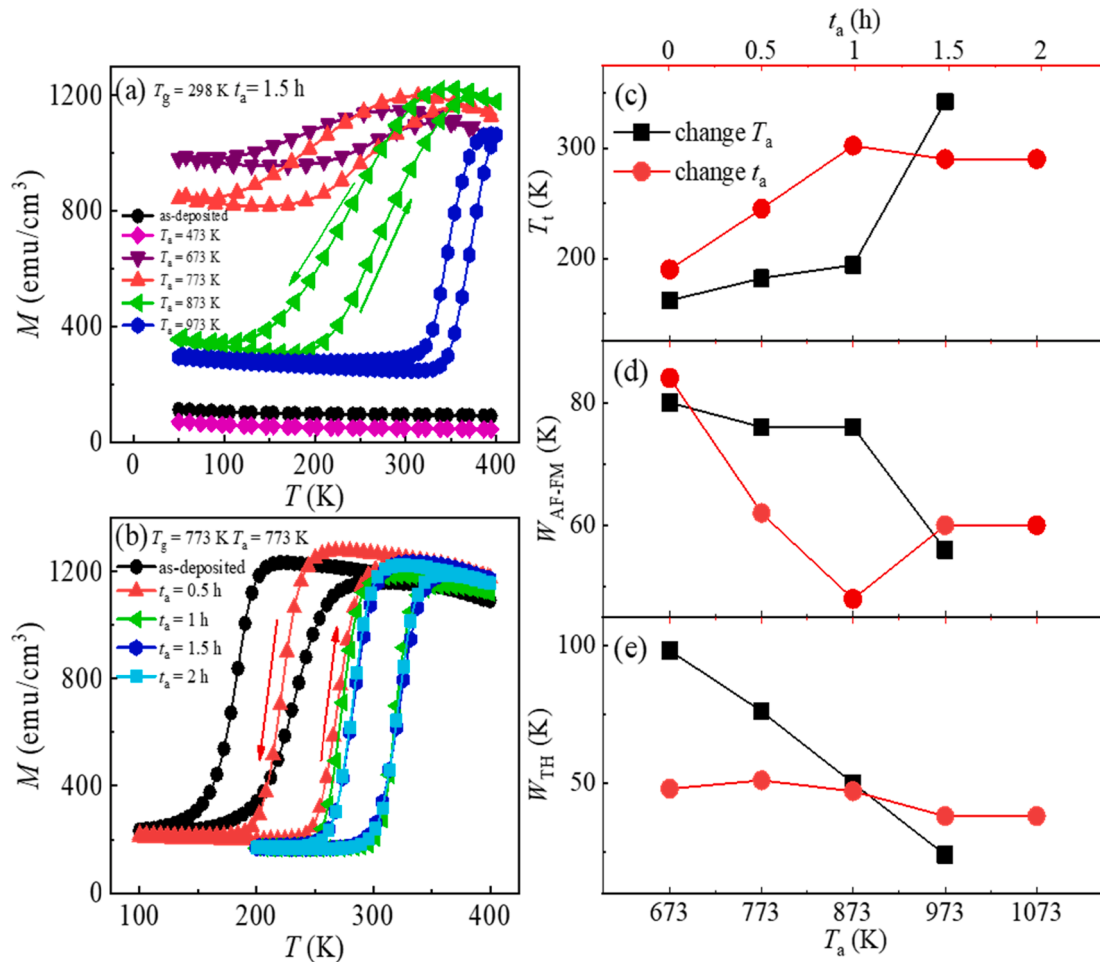


Fig. 3. (a) Temperature-dependence magnetization curves of FeRh films grown at 773 K for different  $T_a$ . (b) Temperature-dependence resistivity  $\rho$  curves of FeRh films for different  $T_a$ . The insets show relationship between  $\rho$  ( $T_{\text{test}} = 400$  K) and  $T_a$  of FeRh thin films. In FeRh films, the relationship between  $T_a$  and magnetic phase transition parameters: (c)  $T_t$ , (d)  $W_{\text{AF-FM}}$  and (e)  $W_{\text{TH}}$ .



**Fig. 4.** (a) Temperature-dependent magnetization curves of FeRh films grown at 298 K for different  $T_a$  (same films). (b) Temperature-dependence magnetization curves of FeRh films grown at 773 K for different  $t_a$ . (c), (d) and (e) are the plots of the magnetic phase transition parameters  $T_t$ ,  $W_{AF-FM}$  and  $W_{TH}$  with change  $T_a$ , respectively. Where the red curve is from the  $M-T$  curve for change  $T_a$  and the black curve is from the  $M-T$  curve for changing  $t_a$ .

increases significantly. It covers the common first-order metamagnetic phase transition interval of FeRh thin film. Fig. 1(c)-(e) show the roughness of the films grown at 773, 873, 1073 K, respectively. The topography results show that when the  $T_g$  of the FeRh thin film reaches 873 K, there are precipitates phase on the surface of the film [26]. The results shown in Fig. 1(f) indicate that for the thin film with  $T_g$  of 773 K, the  $T_t$  is 210 K, which can be increased to 344 K by increasing the  $T_g$  to 1073 K. Meanwhile, the  $W_{AF-FM}$  and  $W_{TH}$  decreases with increasing  $T_g$  as shown in Fig. 1(g)-(h). But the roughness of the thin film surface will also rise sharply. Considering that annealing can increase the order of FeRh B2 phase [27,28], we subsequently performed several annealing processes to control the magnetic phase transition parameters in epitaxial FeRh films.

Fig. 2(a)-(d) show the topography of the film grown at  $T_g = 773$  K, then annealed at different  $T_a$  for 1.5 h. The surface roughness of FeRh films annealed at  $T_a = 773-973$  K is about 0.23 nm (Fig. 2(a)-(c)). It can be seen that there are many precipitates on the FeRh thin film annealed at 1073 K (Fig. 2(d)). Both the temperature dependent magnetization and resistivity curves show that the  $T_a$  can control the magnetic phase transition parameters of FeRh films (Fig. 3(a)-(b)). In contrast to prior studies [29], it shows that the correlation between the change of magnetic phase transition parameters and the morphology is not strong. It shows that the  $T_a$  increases the  $T_t$  of the FeRh films. But when the  $T_a$  reaches 1073 K, the appearance of the precipitates increases the roughness [26].

These films all have a remanence of approximately 200 emu/cm<sup>3</sup>

when they reach the steady state of AF as the temperature decreases. The possible origin of magnetization in the AF phase at room temperature has been proposed, such as substrate strain, changing composition, symmetry breaking, or topography of FeRh films [30-32]. Fig. 3(c)-(e) show the  $T_t$ ,  $W_{AF-FM}$  and  $W_{TH}$  of magnetic phase transition parameters of FeRh films at different  $T_a$ . When  $T_a$  increases,  $T_t$  increases from 190 to 342 K in the magnetization-temperature ( $M-T$ ) curve and from 204 to 344 K in the resistivity-temperature ( $\rho-T$ ) curve, respectively. Meanwhile, both  $W_{AF-FM}$  and  $W_{TH}$  reduce with increasing of  $T_a$ . The magnetic phase transition parameters of FeRh thin film annealed at 1073 K is close to that in bulk FeRh [33].

Fig. 4(a) shows the  $M-T$  curves of FeRh film grown at room temperature, then continuously annealed at different  $T_a$ . Lommel [34] and Ohtani [35] reported that the as-deposited and annealed FeRh films with a  $T_a$  lower than 473 K are not magnetic due to their fcc structure. In the  $T_a$  range of 673-773 K, the film exhibits higher FM, although the magnetization temperature dependence curve shows that it appears the first-order magnetic phase transition of AF-FM. However, when  $T_a$  reaches 873 K, obvious signs of AF phase have been observed at lower temperature. As the  $T_a$  further increased to 973 K, the fraction of FM phase in the film continues to decrease. The  $T_t$ ,  $W_{AF-FM}$  and  $W_{TH}$  of FeRh film grown at room temperature versus  $T_a$  (Fig. 4(c)-(e)) are similar to those of film grown at 773 K (Fig. 3(c)-(e)). In addition, we have explored the influence of  $t_a$  on magnetic phase transition parameters of FeRh film ( $T_g = T_a = 773$  K) as shown in Fig. 4(b). The  $T_t$ ,  $W_{AF-FM}$  and  $W_{TH}$  are also changed with increasing of  $t_a$  (Fig. 4(c)-(e)). The results



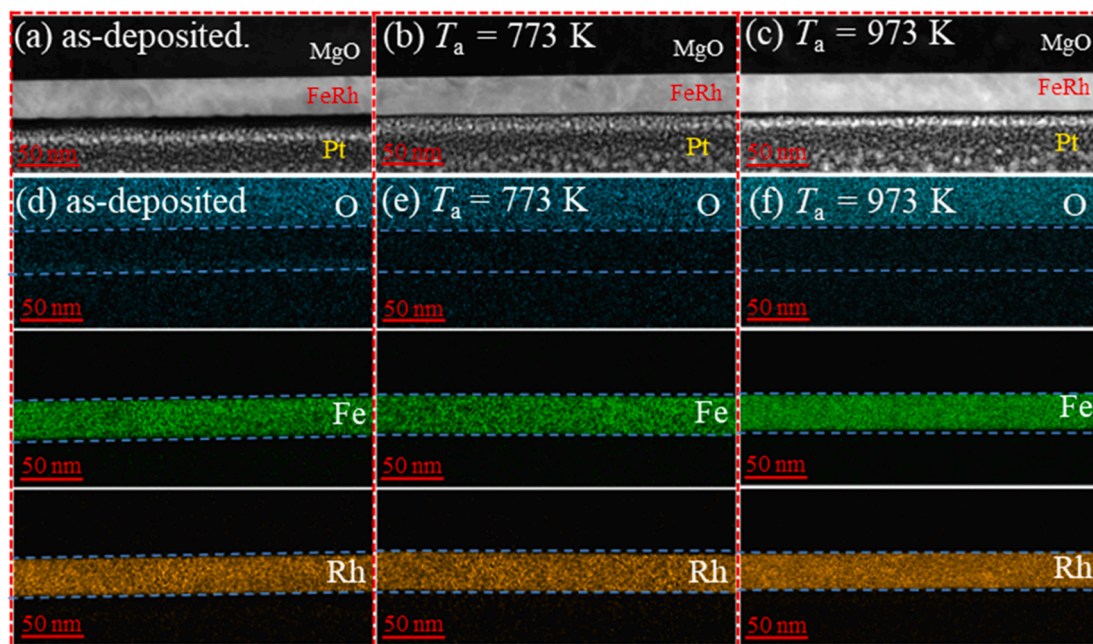


Fig. 5. (a) as-deposited, (b)  $T_a = 773$  K and (c)  $T_a = 973$  K HAADF-STEM images of a cross-sectional TEM lamella showing the FeRh thin film (middle of each image) grown on the MgO substrate (upper), with an average thickness of 30 nm, and protective Pt layer (lower). (d) as-deposited sample, (e)  $T_a = 773$  K and (f)  $T_a = 973$  K chemical composition mapping for O, Fe and Rh obtained using mapping.

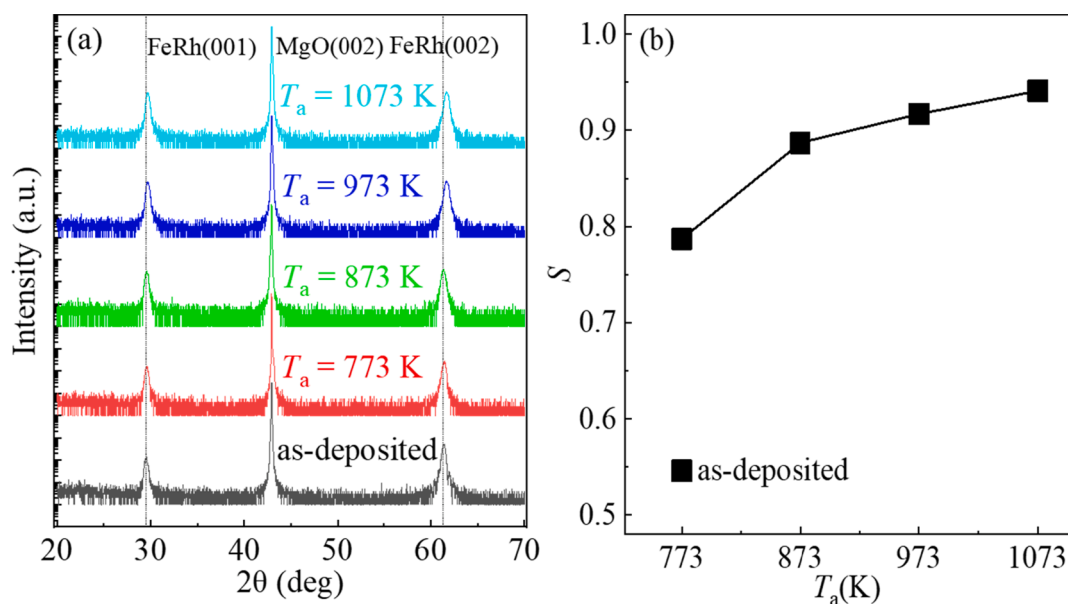
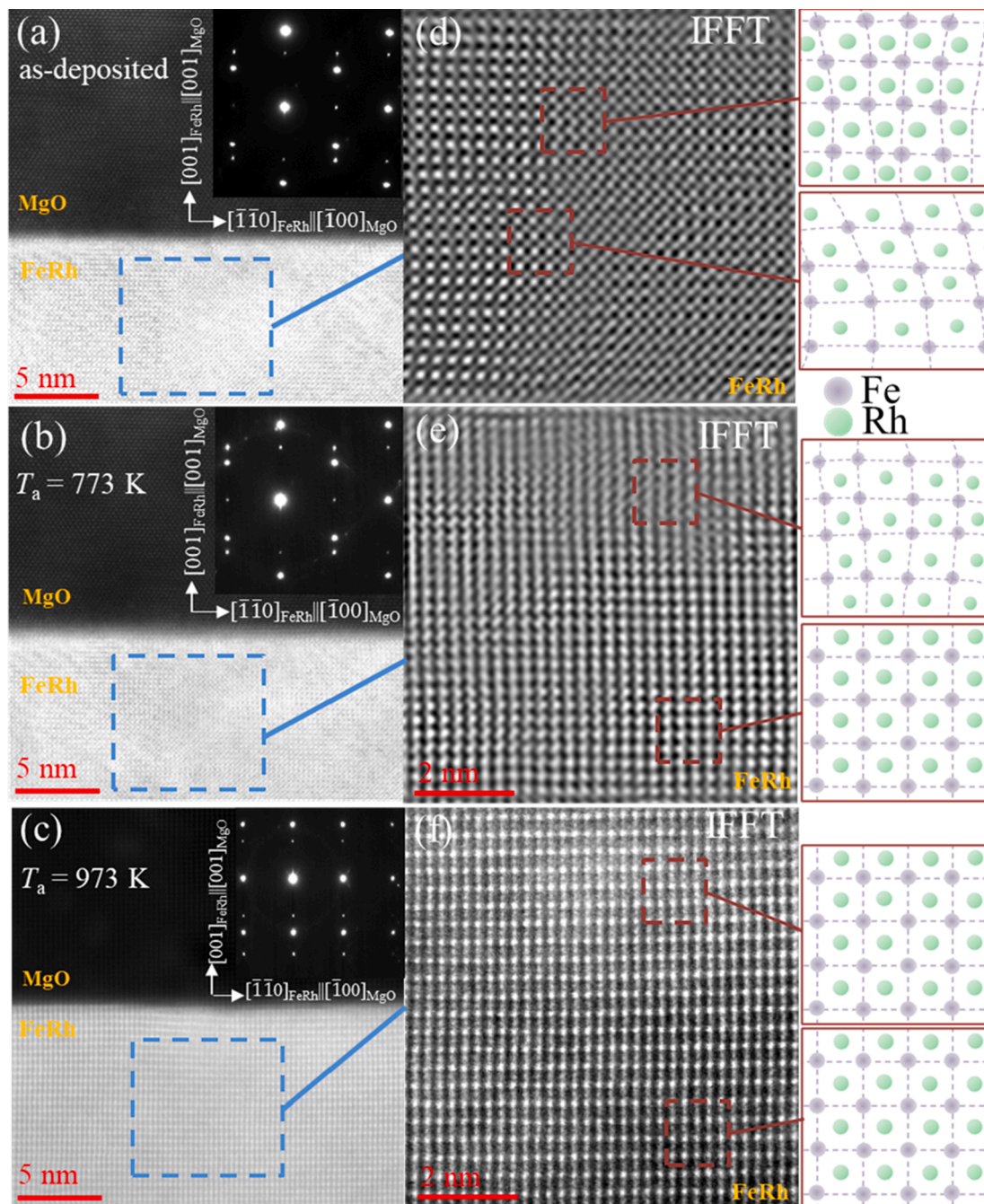


Fig. 6. (a)  $\theta$ - $2\theta$  XRD diffraction patterns of FeRh films for different  $T_a$ . (b) Relationship between the order parameter  $S$  and  $T_a$  of thin films.

show that for the film with  $t_a$  no longer than 1.5 h, the magnetic phase transition parameters are greatly changed with increasing of  $t_a$ . When the  $t_a$  exceeds 1.5 h, the adjustment is weak. All these results indicate that the thermal treatments can regulate the magnetic phase transition parameters of FeRh films.

Fig. 5(a)-(c) show the images of MgO/FeRh/Pt cross sections of as-deposited,  $T_a = 773$  K and  $T_a = 973$  K films in HAADF-STEM mode, respectively. The FeRh films are continuous with uniform growth, and the thickness of the films are about 30 nm. Fig. 5(d)-(f) depict the element composition distribution of O, Fe and Rh of three thermal treatment films, respectively. Neither Fe nor Rh clusters are found in the micro scale. Thus, the thermal treatments may just change the arrangement order degree of Fe (Rh) atoms.

In order to investigate the atomic order degree in FeRh films, we measured the XRD of FeRh films annealed at different  $T_a$  (Fig. 6(a)). The (001) and (002) peaks were found to shift to larger angles, which indicates that the FeRh lattice shrinks. The order parameter ( $S$ ) is a measure of the order quality of the epitaxial FeRh thin film, which is defined as the fraction of the Fe (Rh) atoms occupying the ideal position of the crystal lattice in the B2-FeRh crystal structure of the order condition. According to the mirror XRD scan,  $S$  can be determined as  $S = (I_{(001)}/I_{(002)})^{1/2} / (I_{(001)}^*/I_{(002)}^*)^{1/2}$ , where  $I_{(001)}$  and  $I_{(002)}$  are the integrated intensities of the superlattice Bragg peak and the basic Bragg peak respectively.  $I_{(001)}^*$  and  $I_{(002)}^*$  refer to the theoretical intensities calculated for each sample [36]. As shown in Fig. 6(b), the order parameter  $S$  of FeRh films rises from 0.54 in the as-deposited to 0.92



**Fig. 7.** (a) as-deposited, (b)  $T_a = 773$  K and (c)  $T_a = 973$  K HAADF-STEM images of a cross-sectional atomic scale resolution showing the FeRh films (upper) grown on the MgO (001) substrate (lower). The inset shows the a corresponding SAED pattern. (d), (e) and (f) the a corresponding IFFT pattern. The image on the right corresponds to the partial enlarged schematic diagram of IFFT.

when annealed at 973 K. The change rate of order parameter  $S$  is about 70%. This result indicates that the atomic order degree of the FeRh thin film is positively correlated with the  $T_a$ . This is consistent with the resistivity of thin films annealed at different  $T_a$ . As shown in the inset of Fig. 3(b), at 400 K, as the  $T_a$  increases, the resistivity of thin films decreases with increasing  $T_a$ . In an intact crystal, electrons basically move in a uniform electric field, while in a defective crystal, the periodicity of the lattice in the defect area is destroyed and the electric field changes sharply, which strongly scatters the electrons, resulting in the increase of resistivity in a crystal [37]. Therefore, the decrease of resistivity with the increase of  $T_a$  is due to the tendency of FeRh crystal structure to be ordered [38].

To further illustrate that the  $T_a$  increases the atomic order of FeRh

films, TEM observations were carried out on samples prepared by FIB. Fig. 7(a)-(c) show the atomic images of as-deposited, 773 K and 973 K annealed films, respectively. It shows that FeRh films grows epitaxially on MgO substrate, although there are some defects. The inset of each figure shows the corresponding selected area electron diffraction (SAED) pattern, indicating that the film is a completely dense, epitaxial FeRh layer. It has a B2-ClCs structure, and the film lattice has a  $45^\circ$  in-plane rotation relative to the MgO (001) substrate below:  $(001)_{\text{FeRh}} \parallel (001)_{\text{MgO}}$  and  $[1\bar{1}0]_{\text{FeRh}} \parallel [100]_{\text{MgO}}$  [11]. Fig. 7 (d)-(f) corresponds to inverse fast Fourier transform (IFFT) patterns, showing the projection of the Fe (Rh) atomic column of the corresponding film. The picture on the right is an enlarged schematic diagram of the local area in the corresponding IFFT picture. It shows that there are a large number of atomic



displacements in the as-deposited FeRh film. The atomic displacements in the FeRh film annealed at 773 K are significantly reduced, and almost no atomic displacements are seen in the film annealed at 973 K. Therefore, the increase of  $T_a$  can decrease the displacement of Fe (Rh) atoms in the film to certain extent, then reduces the compression and expansion of the lattice, which may enhance the exchange coupling force of Fe atoms in AF state [39]. As a result,  $T_a$  improves the stability of the FeRh film in the AF state, thereby increasing the  $T_I$  of the film. In the meanwhile, both  $W_{AF-FM}$  and  $W_{TH}$  become narrower, this could be explained by the fact that the increase of  $T_a$  causes most of the atoms in the film to be ordered and the degree of order is close [40,41], which results in the lattice having similar transition temperatures.

#### 4. Conclusions

In summary, by adjusting the  $T_g$  of FeRh to obtain a thin film epitaxially grown in the direction of MgO (001), with AF to FM first-order magnetic phase transition, but when the  $T_g$  reaches 873 K or more, precipitated phases will appear. We have shown that for FeRh films grown at 773 K, the  $T_a$  can increase the AF to FM initial phase transition point from 190 to 342 K. Meanwhile, both  $W_{AF-FM}$  and  $W_{TH}$  are reduced. When the  $T_a$  reaches 1073 K, precipitated phases will also appear. It is worth noting that the experimental results show that the magnetic phase transition parameters changed by annealing are not strongly related to the surface morphology and element distribution of the films. As the  $T_a$  of the FeRh films increases, the perfection of the crystal structure increases, and the atomic order parameter  $S$  varies monotonously between 0.54 and 0.92. In addition, the results prove that annealing is to control the first-order magnetic phase transition by reducing the atomic displacement of thin film. Finally, we want to emphasize that the first-order magnetic phase transition can be controlled by multiple thermal treatments in the same FeRh film. In addition, the first-order magnetic phase transition parameters of FeRh films can be controlled by heat treatments in a wider range, which is essential for the development of several existing and proposed technologies.

#### CRedit authorship contribution statement

**Dezhi Zha:** Investigation, Formal analysis, Writing – original draft. **Baomin Wang:** Conceptualization, Writing – review & editing, Supervision, Funding acquisition. **Lu Yuan:** Investigation. **Yali Xie:** Formal analysis. **Huali Yang:** Formal analysis. **Kai Huang:** Investigation. **Lan Yu:** Supervision, Funding acquisition. **Run-Wei Li:** Supervision, Funding acquisition.

#### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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